

$$p_1(y, z) = p_{11} \cos \pi z, T_1(y, z) = T_{11} \cos \pi z \quad (3.19)$$

where the prime in $v'_{11}(y)$ denote the differentiation with respect to y . Expressions for $v_1(y, z)$ and $w_1(y, z)$ have chosen so that the continuity equation (3.13) is satisfied, substituting these expressions into (3.14) to (3.17), we get

$$\begin{aligned} u''_{11}(y) + R\alpha u'_{11}(y) - \pi^2 u_{11}(y) + \\ \frac{1}{2} Rk_2 (\alpha u''_{11}(y) - \pi^2 \alpha u'_{11}(y) - v_{11}(y) u'''_0 + \\ v''_{11}(y) - \pi^2 v_{11}(y) + 2v'_{11}(y) u''_0) = \\ Rv_{11}(y) u'_0 - R^2 GT_{11} \end{aligned} \quad (3.20)$$

$$\begin{aligned} v''_{11}(y) + R\alpha v'_{11}(y) - (\pi^2 + M^2) v_{11}(y) + \\ \frac{1}{2} Rk_2 \alpha (v''_{11} - \pi^2 v'_{11}) = Rp'_{11}(y) \quad (3.21) \\ v'''_{11}(y) + R\alpha v''_{11}(y) - (\pi^2 + M^2) v'_{11}(y) \\ + \frac{1}{2} Rk_2 \alpha (\pi^2 v''(y) - v'''_{11}(y)) = \pi^2 Rp_{11}(y) \quad (3.22) \end{aligned}$$

$$\begin{aligned} T''_{11}(y) + RP\alpha T'_{11}(y) - \pi^2 T_{11}(y) = \\ RPv_{11}(y) T'_0 \end{aligned} \quad (3.23)$$

Eliminating $P_{11}(y)$ between (3.21) and (3.22), we get

$$\begin{aligned} v'''_{11}(y) + R\alpha v''(y) - (\pi^2 + M^2) \\ v''_{11}(y) - R\pi^2 \alpha v'_{11} + (\pi^4 + \pi^2 M^2) \\ v_{11}(y) = \frac{1}{2} R\alpha k_2 (v'''_{11}(y) - \pi^2 v''(y) + \\ \pi^2 v''_{11} - \pi^4 v'_{11}) \end{aligned} \quad (3.24)$$

The corresponding boundary conditions are

$$\begin{aligned} y=0; u_{11}=0, v_{11}=-\alpha, v'_{11}=0, p_{11}=0, T_{11}=0, \\ y \rightarrow \infty; u_{11}=0, v_{11}=0, v'_{11}=0, p_{11}=0, T_{11}=0 \end{aligned} \quad (3.25)$$

If we consider very small values of viscoelastic parameter k_2 , then substituting

$$u_{11}=u_{110}(y)+k_2 u_{111}(y)+O(k^2_2) \\ v_{11}=v_{110}(y)+k_2 v_{111}(y)+O(k^2_2) \quad (3.26)$$

in the equations (3.20) and (3.24) and equating the coefficient of like powers of k_2 , we get

$$u''_{110} + R\alpha u'_{110} - \pi^2 u_{110} = Rv_{110} u'_0 - R^2 GT_{11} \quad (3.27)$$

$$u''_{111} + R\alpha u'_{111} - \pi^2 u_{111} = \frac{1}{2} R [2v_{110} u'_0 - \\ \alpha u''_{110} + \pi^2 \alpha u'_{110} + v_{110} u'''_{110} - v_{110} + \\ \pi^2 v_{110} - 2v'_{110} u''_0] \quad (3.28)$$

$$\begin{aligned} v'''_{110}(y) + R\alpha v'''_{110}(y) - (2\pi^2 + M^2) \\ v''_{110}(y) - R\pi^2 \alpha v'_{110} + (\pi^4 + \pi^2 M^2) v_{110} = 0 \end{aligned} \quad (3.29)$$

$$\begin{aligned} v'''_{111}(y) + R\alpha v'''_{111}(y) - (2\pi^2 + M^2) \\ v''_{111}(y) - R\pi^2 \alpha v'_{111} + (\pi^4 + \pi^2 M^2) \\ v_{111} = \frac{1}{2} R\alpha [v'''_{110} + \pi^2 v'''_{110} + \pi^2 v''_{110} - \pi^4 v'_{110}] \end{aligned} \quad (3.30)$$

$$\begin{aligned} \text{The corresponding boundary conditions are} \\ y=0; u_{110}=0, u_{111}=0, v_{110}=-\alpha, v_{111}=0 \\ y \rightarrow \infty; u_{110}=0, u_{111}=0, v_{110}=0, v_{111}=0 \end{aligned} \quad (3.31)$$

The differential equations (3.27) to (3.30) can be solved with the help of the boundary conditions (3.31).

The solutions of (3.27) to (3.30) and (3.23) are as follows:

$$\begin{aligned} u_{110} &= B_{29} e^{-\lambda_4 y} + B_{22} e^{-m_1 y} + B_{23} e^{-m_2 y} + B_{24} e^{-m_3 y} \\ B_{25} e^{-m_4 y} + B_{26} e^{-m_5 y} - B_{27} e^{-m_6 y} - B_{28} e^{-\alpha_2 y} \\ u_{111} &= B_{38} e^{-\lambda_4 y} + B_{39} e^{-\alpha_2 y} + B_{40} e^{-m_1 y} \\ + B_{41} e^{-m_2 y} + B_{42} e^{-m_3 y} \end{aligned}$$